# General boundary conditions for the Boussinesq-Abbott model with varying topography

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# Objectives

#### Project supported by IMPT and ANR Bourgeons

Supervision : David Lannes and Philippe Bonneton

Long term goal: study extreme waves in littoral area



Mathieu Rigal

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# Outline

- The model & problematic
- 2 Nonlocal reformulation for flat bottom
- Extention to varying bottom
- Output State Numerical experiment

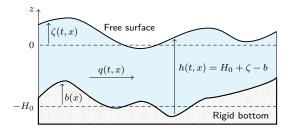
# The PDEs

Boussinesq-Abbott model with varying bottom:

$$\begin{cases} \partial_t \zeta + \partial_x q = 0\\ (1 + h_b \mathcal{T}) \partial_t q + \partial_x f(\zeta, q) = -gh \partial_x b \end{cases}$$

with  $h_b = H_0 - b$  (depth at rest) and

$$f(\zeta, q) = \frac{q^2}{h} + \frac{gh^2}{2}, \qquad h_b \mathcal{T}(\cdot) = -\frac{1}{3}\partial_x(h_b^2\partial_x \cdot) + \text{ lower order terms}$$



(BA)

# The PDEs

Initial-boundary value problem:

$$\begin{cases} \partial_t \zeta + \partial_x q = 0\\ (1 + h_b \mathcal{T}) \partial_t q + \partial_x f(\zeta, q) = -gh \partial_x b \end{cases}$$
 in (0,  $\ell$ ) (BA)

completed with

$$(\zeta, q)_{|_{t=0}} = (\zeta_0, q_0), \qquad \zeta(t, 0) = g_0(t), \qquad \zeta(t, \ell) = g_\ell(t)$$

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#### How to recover $q_{|_{x=0,\ell}}(t)$ ? (missing data)

- Hyperbolic case  $(h_b \mathcal{T} \equiv 0)$ : Riemann invariants
- Dispersive case: need to invert  $(1 + h_b T) \rightarrow$  requires knowledge on  $\partial_t q_{|_{x=0,\ell}}$

[WKS99] Wei, Kirby, Sinha. "Generation of waves in Boussinesq models using a source function method" (1999)

- [DMS09] Dougalis, Mitsotakis, Saut. "On initial-boundary value problems for a Boussinesq system of BBM-BBM type in a plane domain" (2009)
- [LW20] Lannes, Weynans. "Generating boundary conditions for a Boussinesq system" (2020)
- [BL22] Beck, Lannes. "Freely floating objects on a fluid governed by the Boussinesq equations" (2022)
- [LR24] Lannes, R. "General boundary conditions for a Boussinesq model with varying bathymetry" (2024)

Flat bottom case ( $b \equiv Cst$ ): discharge eq. simplifies to

$$(1 - \kappa^2 \partial_{xx}^2) \partial_t q + \partial_x f(\zeta, q) = 0 \qquad \text{in } (0, \ell)$$

Fix  $0 \le t \le T$ , then  $y(x) = \partial_t q(t, x)$  satisfies an ODE of the form  $\begin{cases} y - \kappa^2 y'' = \phi(x) \\ y(0) = \dot{q}_{|_0}, \quad y(\ell) = \dot{q}_{|_\ell} \end{cases}$ Equivalently:  $y = y_h + y_b$  with  $\begin{cases} y_h - \kappa^2 y''_h = 0 \\ y_h(0) = \dot{q}_{|_0}, \quad y_h(\ell) = \dot{q}_{|_\ell} \end{cases}$  and  $\begin{cases} y_b - \kappa^2 y''_b = \phi(x) \\ y_b(0) = y_b(\ell) = 0 \end{cases}$ 

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Define  $R^0$  as the inverse of  $(1 - \kappa^2 \partial_{xx}^2)$  with homogeneous Dirichlet conditions at  $x = 0, \ell$ 

$$\Rightarrow \partial_{t}q = \underbrace{-R^{0}\partial_{x}f}_{y_{b}} + \underbrace{\mathfrak{s}_{(0)}(x)\dot{q}_{|_{x=0}} + \mathfrak{s}_{(\ell)}(x)\dot{q}_{|_{x=\ell}}}_{y_{h}}$$
  
where 
$$\begin{cases} (1 - \kappa^{2}\partial_{xx}^{2})\mathfrak{s}_{(0)} = 0\\ \mathfrak{s}_{(0)}(0) = 1, \quad \mathfrak{s}_{(0)}(\ell) = 0 \end{cases} \text{ and } \begin{cases} (1 - \kappa^{2}\partial_{xx}^{2})\mathfrak{s}_{(\ell)} = 0\\ \mathfrak{s}_{(\ell)}(0) = 0, \quad \mathfrak{s}_{(\ell)}(\ell) = 1 \end{cases} .$$
(1)

Note  $R^1$  the inverse of  $(1 - \kappa^2 \partial_{xx}^2)$  with **homogeneous Neumann** conditions at  $x = 0, \ell$ 

 $\Rightarrow R^0 \partial_x = \partial_x R^1$ 

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#### Proposition 1 (D. Lannes, R.)

Let  $(\zeta, q)$  be such that the **compatibility conditions** hold:

$$\begin{cases} \zeta_{|_{t=0}}(0) = g_0(0) \\ -\partial_x q_{|_{t=0}}(0) = \dot{g}_0(0) \end{cases}, \qquad \begin{cases} \zeta_{|_{t=0}}(\ell) = g_\ell(0) \\ -\partial_x q_{|_{t=0}}(\ell) = \dot{g}_\ell(0) \end{cases}$$

Then the two assertions are equivalent:

• The pair ( $\zeta$ , q) satisfies the IBVP (BA) with  $\zeta(\cdot, 0) = g_0$  and  $\zeta(\cdot, \ell) = g_\ell$ 

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2 The pair  $(\zeta, q)$  satisfies the IVP

$$\begin{cases} \partial_t \zeta + \partial_x q = 0\\ \partial_t q + \partial_x (R^1 f) = \mathfrak{s}_{(0)} \dot{q}_{|_0} + \mathfrak{s}_{(\ell)} \dot{q}_{|_\ell} \end{cases} \quad in \ (0, \ell),$$

with the trace equations

$$\begin{pmatrix} \mathbf{s}_{(0)}'(0) & \mathbf{s}_{(\ell)}'(0) \\ \mathbf{s}_{(0)}'(\ell) & \mathbf{s}_{(\ell)}'(\ell) \end{pmatrix} \begin{pmatrix} \dot{\mathbf{q}}_{|_{\ell}} \\ \dot{\mathbf{q}}_{|_{\ell}} \end{pmatrix} = \frac{1}{\kappa^2} \begin{pmatrix} (R^1 - \mathrm{id})_{|_{\ell}} f \\ (R^1 - \mathrm{id})_{|_{\ell}} f \end{pmatrix} - \begin{pmatrix} \ddot{\mathbf{g}}_0 \\ \ddot{\mathbf{g}}_\ell \end{pmatrix}$$

(3)

(2)

Proof:

 $\Rightarrow$ 

$$\begin{aligned} \partial_t \zeta + \partial_x q &= 0 \\ \partial_t q + \partial_x (R^1 f) &= \mathfrak{s}_{(0)} \dot{q}_{|_0} + \mathfrak{s}_{(\ell)} \dot{q}_{|_{\ell}} \end{aligned}$$

Proof:  $\Rightarrow$ 

$$\begin{cases} \partial_t \zeta + \partial_x q = 0\\ \frac{\partial_x (\partial_t q + \partial_x (R^1 f) = \mathfrak{s}_{(0)} \dot{q}_{|_0} + \mathfrak{s}_{(\ell)} \dot{q}_{|_\ell}) \end{cases}$$

$$\Rightarrow \quad -\partial_{tt}\zeta + \partial_{xx}^2(R^1f) = \mathfrak{s}_{(0)}'\dot{q}_{|_0} + \mathfrak{s}_{(\ell)}'\dot{q}_{|_\ell}$$

Proof:  $\Rightarrow$ 

$$\begin{cases} \partial_t \zeta + \partial_x q = 0\\ \partial_x \left( \partial_t q + \partial_x (R^1 f) = \mathfrak{s}_{(0)} \dot{q}_{|_0} + \mathfrak{s}_{(\ell)} \dot{q}_{|_\ell} \right) \qquad \Rightarrow \qquad -\partial_{tt} \zeta + \underbrace{\partial^2_{xx} (R^1 f)}_{\frac{1}{k^2} (R^1 - \mathrm{id})f} = \mathfrak{s}_{(0)}' \dot{q}_{|_0} + \mathfrak{s}_{(\ell)}' \dot{q}_{|_\ell} \end{cases}$$

By construction

$$(1 - \kappa^2 \partial_{xx}^2) R^1 = \mathrm{id} \implies \partial_{xx}^2 R^1 = \kappa^{-2} (R^1 - \mathrm{id})$$

Proof:

$$\begin{cases} \partial_t \zeta + \partial_x q = 0\\ \partial_x \left( \partial_t q + \partial_x (R^1 f) = \mathfrak{s}_{(0)} \dot{q}_{|_0} + \mathfrak{s}_{(\ell)} \dot{q}_{|_\ell} \right) \qquad \Rightarrow \qquad -\partial_{tt} \zeta + \underbrace{\partial^2_{xx} (R^1 f)}_{\frac{1}{k^2} (R^1 - \mathrm{id})f} = \mathfrak{s}_{(0)}' \dot{q}_{|_0} + \mathfrak{s}_{(\ell)}' \dot{q}_{|_\ell} \end{cases}$$

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Taking the trace at x = 0,  $\ell$  and using that  $\zeta_{|_{x=0,\ell}} = g_{0,\ell}$ :

$$\begin{cases} -\ddot{g}_{0} + \frac{1}{\kappa^{2}}(R^{1} - \mathrm{id})_{|_{x=0}}f(\zeta, q) = \mathfrak{s}'_{(0)}(0)\dot{q}_{|_{0}} + \mathfrak{s}'_{(\ell)}(0)\dot{q}_{|_{\ell}} \\ -\ddot{g}_{\ell} + \frac{1}{\kappa^{2}}(R^{1} - \mathrm{id})_{|_{x=\ell}}f(\zeta, q) = \mathfrak{s}'_{(0)}(\ell)\dot{q}_{|_{0}} + \mathfrak{s}'_{(\ell)}(\ell)\dot{q}_{|_{\ell}} \end{cases}$$

**Proof:** 
$$\leftarrow$$
 We now assume that  

$$\begin{cases}
-\ddot{g}_{0} + \frac{1}{\kappa^{2}}(R^{1} - \mathrm{id})_{|_{x=0}}f(\zeta, q) = s'_{(0)}(0)\dot{q}_{|_{0}} + s'_{(\ell)}(0)\dot{q}_{|_{\ell}} \\
-\ddot{g}_{\ell} + \frac{1}{\kappa^{2}}(R^{1} - \mathrm{id})_{|_{x=\ell}}f(\zeta, q) = s'_{(0)}(\ell)\dot{q}_{|_{0}} + s'_{(\ell)}(\ell)\dot{q}_{|_{\ell}}
\end{cases}$$
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$$\& \qquad \begin{cases} \zeta_{|_{t=0}}(x=0,\ell) = g_{0,\ell}(0) \\ -\partial_x q_{|_{t=0}}(x=0,\ell) = \dot{g}_{0,\ell}(0) \end{cases}$$

Following previous steps we also have

$$\begin{cases} -\ddot{\zeta}_{|_{0}} + \frac{1}{\kappa^{2}}(R^{1} - \mathrm{id})_{|_{x=0}}f(\zeta, q) = \mathfrak{s}_{(0)}'(0)\dot{q}_{|_{0}} + \mathfrak{s}_{(\ell)}'(0)\dot{q}_{|_{\ell}} \\ -\ddot{\zeta}_{|_{\ell}} + \frac{1}{\kappa^{2}}(R^{1} - \mathrm{id})_{|_{x=\ell}}f(\zeta, q) = \mathfrak{s}_{(0)}'(\ell)\dot{q}_{|_{0}} + \mathfrak{s}_{(\ell)}'(\ell)\dot{q}_{|_{\ell}} \end{cases}$$

$$\begin{array}{ll} \textbf{Proof:} & \overleftarrow{\leftarrow} & \text{We now assume that} \\ \\ & \left\{ \begin{split} -\ddot{g}_{0} + \frac{1}{\kappa^{2}} (R^{1} - \text{id})_{|_{x=0}} f(\zeta, q) &= s_{(0)}'(0) \dot{q}_{|_{0}} + s_{(\ell)}'(0) \dot{q}_{|_{\ell}} \\ - \ddot{g}_{\ell} + \frac{1}{\kappa^{2}} (R^{1} - \text{id})_{|_{x=\ell}} f(\zeta, q) &= s_{(0)}'(\ell) \dot{q}_{|_{0}} + s_{(\ell)}'(\ell) \dot{q}_{|_{\ell}} \end{split} \right. \qquad & \left\{ \begin{split} \zeta_{|_{t=0}}(x = 0, \ell) &= g_{0,\ell}(0) \\ -\partial_{x} q_{|_{t=0}}(x = 0, \ell) &= \dot{g}_{0,\ell}(0) \end{split} \right. \end{aligned}$$

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$$\begin{cases} -\ddot{\zeta}_{|_{0}} + \frac{1}{\kappa^{2}}(R^{1} - \mathrm{id})_{|_{x=0}}f(\zeta, q) = s'_{(0)}(0)\dot{q}_{|_{0}} + s'_{(\ell)}(0)\dot{q}_{|_{\ell}} \\ \\ -\ddot{\zeta}_{|_{\ell}} + \frac{1}{\kappa^{2}}(R^{1} - \mathrm{id})_{|_{x=\ell}}f(\zeta, q) = s'_{(0)}(\ell)\dot{q}_{|_{0}} + s'_{(\ell)}(\ell)\dot{q}_{|_{\ell}} \end{cases} \Rightarrow \begin{cases} \ddot{\zeta}_{|_{0}} = \ddot{g}_{0} \\ \\ \ddot{\zeta}_{|_{\ell}} = \ddot{g}_{\ell} \end{cases}$$

**Compatibility conditions**  $\Rightarrow$  solution given by  $(\zeta_{i_0}, \zeta_{i_\ell}) = (g_0, g_\ell)$ 

Possibility to enforce general boundary conditions  $\xi_{0}^{+}(\zeta_{|_{0}}, q_{|_{0}})(t) = g_{0}(t), \qquad \xi_{\ell}^{-}(\zeta_{|_{\ell}}, q_{|_{\ell}})(t) = g_{\ell}(t). \tag{4}$ For instance,  $\xi^{\pm}$  given by q or Saint-Venant Riemann invariants  $\mathcal{R}_{\pm}(U) = \frac{q}{h} \pm 2\sqrt{gh}$  Possibility to enforce general boundary conditions  $\xi_0^+(\zeta_0, q_0)(t) = g_0(t), \qquad \xi_\ell^-(\zeta_\ell, q_\ell)(t) = g_\ell(t). \tag{4}$ For instance,  $\xi^{\pm}$  given by q or Saint-Venant Riemann invariants  $\mathcal{R}_{\pm}(U) = \frac{q}{h} \pm 2\sqrt{gh}$ 

Adapt trace ODE in terms of missing data (outgoing information  $\xi_0^-$  and  $\xi_\ell^+$ )

$$\begin{cases} s'_{(0)}(0) & s'_{(\ell)}(0) \\ s'_{(0)}(\ell) & s'_{(\ell)}(\ell) \end{cases} \frac{d}{dt} \begin{pmatrix} q(\xi_0^+, \xi_0^-) \\ q(\xi_\ell^+, \xi_\ell^-) \end{pmatrix} = \frac{1}{\kappa^2} \begin{pmatrix} (R^1 - id)_{|_0} f \\ (R^1 - id)_{|_\ell} f \end{pmatrix} - \frac{d^2}{dt^2} \begin{pmatrix} \zeta(\xi_0^+, \xi_0^-) \\ \zeta(\xi_\ell^+, \xi_\ell^-) \end{pmatrix}$$

# Case of a varying topography ( $b \neq Cst$ )

$$(1+h_b\mathcal{T})\partial_t q+\partial_x f(\zeta,q)=-gh\partial_x b$$

where  $h_b \mathcal{T}(\cdot) = -\frac{1}{3} \partial_x (h_b^2 \partial_x \cdot) + \text{ lower order terms}$ 

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Note  $R_b^0$  the inverse of  $(1 + h_b T)$  with homogeneous Dirichlet cond. and

$$\begin{cases} (1+h_b\mathcal{T})\mathfrak{s}_{(b,0)} = 0\\ \mathfrak{s}_{(b,0)}(0) = 1, \quad \mathfrak{s}_{(b,0)}(\ell) = 0 \end{cases} \begin{cases} (1+h_b\mathcal{T})\mathfrak{s}_{(b,\ell)} = 0\\ \mathfrak{s}_{(b,\ell)}(0) = 0, \quad \mathfrak{s}_{(b,\ell)}(\ell) = 1 \end{cases}$$

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$$\Rightarrow \left| \partial_t q = -R_b^0 \partial_x f - g R_b^0 (h \partial_x b) + \mathfrak{s}_{(b,0)} \dot{q}_{|_0} + \mathfrak{s}_{(b,\ell)} \dot{q}_{|_\ell} \right|$$

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Next step: generalisation of  $R^0 \partial_x = \partial_x R^1$ ?

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However we can construct  $R_b^1$  verifying for some  $(\alpha, \beta)(x)$ 

$$R_b^0(\partial_x + \alpha) = (\partial_x + \beta)R_b^1$$

$$\Rightarrow \partial_t q = -\partial_x R_b^1 f \underbrace{-\beta R_b^1 f + R_b^0(\alpha f - gh\partial_x b)}_{\mathfrak{B}(\zeta, q)} + \mathfrak{s}_{(b,0)} \dot{q}_{|_0} + \mathfrak{s}_{(b,\ell)} \dot{q}_{|_\ell}$$

#### Proposition 2 (D. Lannes, R.)

Let  $(\zeta, q)$  be such that the **compatibility conditions** hold:

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Then the two assertions are equivalent:

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- 2 The pair  $(\zeta, q)$  satisfies the IVP

$$\begin{cases} \partial_t \zeta + \partial_x q = 0\\ \partial_t q + \partial_x (R_b^1 f) = \mathfrak{B}(\zeta, q) + \mathfrak{s}_{(b,0)} \dot{q}_{|_0} + \mathfrak{s}_{(b,\ell)} \dot{q}_{|_\ell} \end{cases} \quad in \ (0,\ell), \tag{5}$$

with the trace equations

$$\begin{pmatrix} \varsigma_{(b,0)}^{\prime}(0) & \varsigma_{(b,\ell)}^{\prime}(0) \\ \varsigma_{(b,0)}^{\prime}(\ell) & \varsigma_{(b,\ell)}^{\prime}(\ell) \end{pmatrix} \begin{pmatrix} \dot{q}_{|_{\ell}} \\ \dot{q}_{|_{\ell}} \end{pmatrix} + V_{boundary}(\zeta_{|_{0,\ell}}, q_{|_{0,\ell}}) = V_{interior}[\zeta, q] - \begin{pmatrix} \ddot{g}_{0} \\ \ddot{g}_{\ell} \end{pmatrix}$$
(6)

where V<sub>boundary</sub>, V<sub>interior</sub> are known.

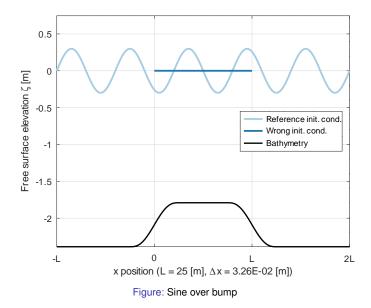
- Standard finite differences for nonlocal terms and trace equations
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#### Asymptotic stability

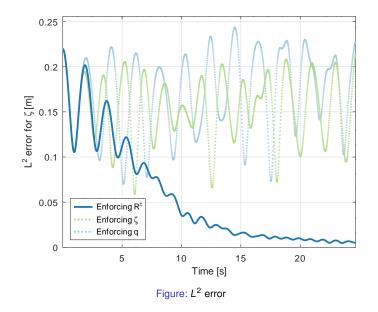
**Question:** starting from a wrong initial condition, can we recover the reference solution by enforcing appropriate boundary conditions?

# Numerical experiment



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2

# Conclusion

- Boussinesq-Abbott model is accurate, but boundary conditions are challenging
- Reformulation strategy allows to recover missing data
- Extension to varying bathymetries & general boundary conditions
- Numerical validation and experiments (asymptotic stability)

Perspectives: high order DGFEM code, improved dispersion relation

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Thank you for your attention!