

# Quasi-static Free-Boundary Equilibrium of Toroidal Plasma: Computational Methods and Applications

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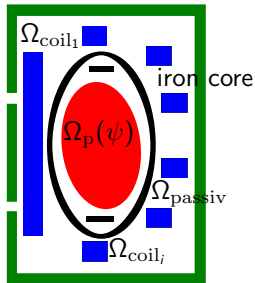
# Quasi-static Free-Boundary Equilibrium of Toroidal Plasma

inside plasma and non-conducting parts:

$$\mathbf{grad} p = \mathbf{J} \times \mathbf{B}, \quad \mathbf{div} \mathbf{B} = 0, \quad \mathbf{curl} \frac{1}{\mu} \mathbf{B} = \mathbf{J},$$

in all other conducting structures:

$$\partial_t \mathbf{B} = \mathbf{curl} \frac{1}{\sigma} (\mathbf{J}_{src} - \mathbf{J}), \quad \mathbf{div} \mathbf{B} = 0, \quad \mathbf{curl} \frac{1}{\mu} \mathbf{B} = \mathbf{J}.$$



with toroidal symmetry: ( $\psi$  toroidal comp. of  $r \mathbf{A}$ ,  $\mathbf{B} = \mathbf{curl} \mathbf{A}$ )

$$-\nabla \left( \frac{1}{\mu[\psi]r} \nabla \psi \right) = \begin{cases} rp'(\psi) + \frac{1}{\mu_0 r} ff'(\psi) & \text{in } \Omega_P(\psi), \\ \frac{n_i V_i(t)}{R_i S_i} - 2\pi \frac{n_i^2}{R_i S_i^2} \int_{\Omega_{coil_i}} \frac{\partial \psi}{\partial t} dr dz =: \frac{l_i(\psi)}{S_i} & \text{in } \Omega_{coil_i}, \\ -\frac{\sigma}{r} \frac{\partial \psi}{\partial t} & \text{in } \Omega_{passive}, \\ 0 & \text{elsewhere,} \end{cases}$$

with  $p'$  and  $ff'$  known.  $f$  toroidal component of  $r \mathbf{B}$ .

Infinite domain, plasma domain  $\Omega_P(\psi)$  unknown, circuit equations, iron core.

# Genealogy

FEM for free-boundary equilibrium in tor. symm.:

Challenges	SCED [BFT81]	Proteus [ABB87]	CEDRES++ [G99]
iron core	+++	+++	+++[Boulbe '11]
infinite domain	not	+++ [ABB86]	+++ [G99]
free boundary	+++	+++	+++
circuit equations	version Blum	version Albanese	Blum/Boulbe/Nardon '14
<b>Newton</b>	+++	+++	+++ [Hetal15]
inverse problem	static [B89]	stat.	stat. & dynam. [Hetal15]

[BFT81] J. Blum, J. Le Foll, B. Thooris, *The self-consistent equilibrium and diffusion code SCED*, CPC, 1981.

[ABB86] R. Albanese, J. Blum, O. Barbieri, *On the solution of the magnetic flux equation in an infinite domain*, 8th Europhys. Conf. Comp. in Plasma Phys., 1986.

[ABB87] R. Albanese, J. Blum, O. Barbieri, 12th Conf. on Num. Simul. of Plasma, 1987.

[B89] J. Blum, *Numerical simulation and optimal control in plasma physics*, 1989.

[G99] V. Grandgirard, *Modelisation de l'equilibre d'un plasma de tokamak*, PhD., 1999.

[Hetal15] H.H. et al., *Quasi-static free-boundary equilibrium of toroidal plasma with CEDRES++ ...*, Journal of Plasma Physics 2015.

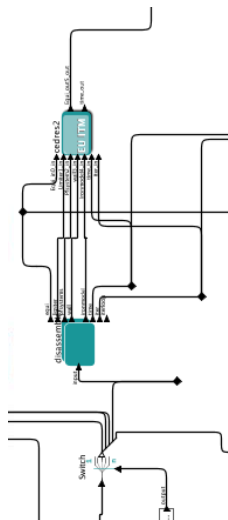
# CEDRES++ & FEEQS.M (C. Boulbe, B. Faugeras, H. H.)

Applications, **focus on control and scenarios:**

- ▶ static equilibrium calculations for given currents,
- ▶ currents calculation for given static equilibrium,
- ▶ evolution of equilibrium calculation for given voltages,
- ▶ voltage evolution calculation for given equil. evolution,
- ▶ **not** real-time reconstruction, **not** for MHD instability

Codes at CASTOR/CEA/UNS

- ▶ CEDRES++
  - ▶ *Couplage Equilibre Diffusion Resistive pour l'Étude des Scenarios*
  - ▶ productive code written in C++;
  - ▶ in use at CEA for set up experiment scenarios at upcoming WEST;
  - ▶ interfaced for ITM (integrated tokamak modeling) Kepler Platform;
- ▶ FEEQS.M



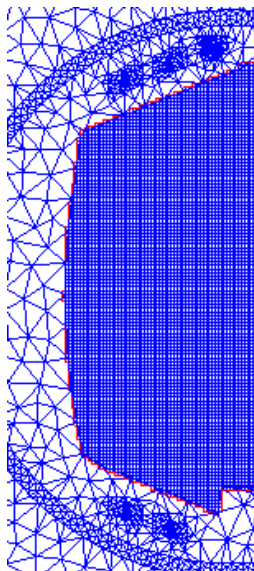
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Codes at CASTOR (C. Boulbe, B. Faugeras, H. H.)

- ▶ CEDRES++
- ▶ FEEQS.M
  - ▶ *Finite Element Equilibrium Solver in Matlab*
  - ▶ same core functionality as CEDRES++;
  - ▶ test and fast prototyping environment in MATLAB;
  - ▶ high performance, thanks to vectorization;
  - ▶ new applications;
  - ▶ improve and extend numerical methods;
  - ▶ simple code distributions;
  - ▶ intern and master projects;



# Outline

## Quasi-Static Free-Boundary Equilibrium of Toroidal Plasma

- Direct Static Problem

- Inverse Static Problem

- Direct Evolution Problem

- Inverse Evolution Problem

## Weak Formulation

## Newton's Method

## Sequential Quadratic Programming

## Validation & Performance

## Application: Vertical Displacement

## Conclusions & Outlook

# What's next?

## Quasi-Static Free-Boundary Equilibrium of Toroidal Plasma

Direct Static Problem

Inverse Static Problem

Direct Evolution Problem

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## Direct Static Problem (for flux $\psi(r, z)$ )

$$-\nabla \cdot \left( \frac{1}{\mu r} \nabla \psi \right) = \begin{cases} r p'(\psi) + \frac{1}{\mu_0 r} f f'(\psi) & \text{in } \Omega_p(\psi), \\ \frac{I_i}{S_i} & \text{in } \Omega_{\text{coil}_i}, \\ 0 & \text{elsewhere,} \end{cases}$$

$$\psi(0, z) = 0, \quad \lim_{\|(r,z)\| \rightarrow +\infty} \psi(r, z) = 0,$$

Iron ( $\mu_{\text{Fe}}$  experimental data):

$$\mu = \mu(r, |\nabla \psi|) = \begin{cases} \mu_{\text{Fe}} (|\nabla \psi|^2 r^{-2}) & \text{in } \Omega_{\text{iron}}, \\ \mu_0 & \text{elsewhere.} \end{cases}$$

Model for current density ( $\alpha, \beta, \gamma, r_0$  given):

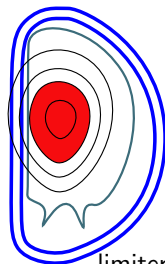
$$p'(\psi) \approx S_{p'}(\psi_N) = \frac{\beta}{r_0} (1 - \psi_N^\alpha)^\gamma,$$

$$f f'(\psi) \approx S_{f f'}(\psi_N) = (1 - \beta) \mu_0 r_0 (1 - \psi_N^\alpha)^\gamma,$$

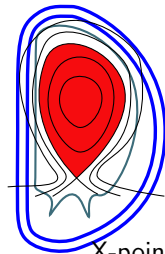
$$\psi_N(r, z) = \frac{\psi(r, z) - \psi_{\text{ax}}[\psi]}{\psi_{\text{bd}}[\psi] - \psi_{\text{ax}}[\psi]},$$

$$\psi_{\text{ax}}[\psi] := \psi(r_{\text{ax}}[\psi], z_{\text{ax}}[\psi]),$$

$$\psi_{\text{bd}}[\psi] := \psi(r_{\text{bd}}[\psi], z_{\text{bd}}[\psi]).$$



limiter plasma



X-point plasma



# Inverse Static Problem (for currents $I_j$ )

Objective and Regularization

$$K(\psi) := \frac{1}{2} \sum_{i=1}^{N_{\text{desi}}} (\psi(r_i, z_i) - \psi(r_{\text{desi}}, z_{\text{desi}}))^2, \quad R(I_1, \dots, I_L) := \sum_{i=1}^L \frac{w_i}{2} I_i^2$$

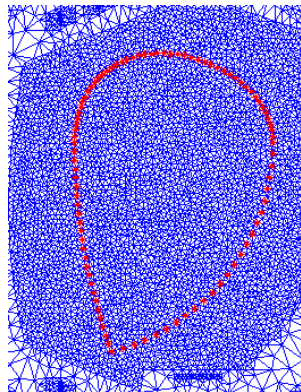
Optimal Control/Inverse Problem:

$$\min_{\psi, I_1, \dots, I_L} K(\psi) + R(I_1, \dots, I_L)$$

subject to

$$-\nabla \cdot \left( \frac{1}{\mu r} \nabla \psi \right) = \begin{cases} r S'_p(\psi_N) + \frac{1}{\mu_0 r} S_{ff'}(\psi_N) & \text{in } \Omega_P(\psi), \\ \frac{I_i}{S_i} & \text{in } \Omega_{\text{coil}_i}, \\ 0 & \text{elsewhere,} \end{cases}$$

$$\psi(0, z) = 0, \quad \lim_{\|(r,z)\| \rightarrow +\infty} \psi(r, z) = 0,$$



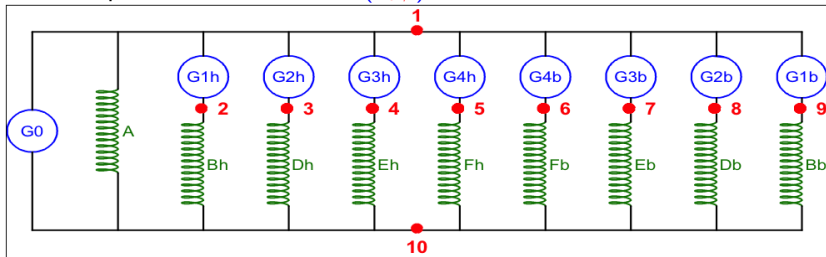
PDE-constrained optimization with **non-linear** constraints.

# Direct Evolution Problem (for flux evolution $\psi(r, z, t)$ )

$$-\nabla \cdot \left( \frac{1}{\mu r} \nabla \psi \right) = \begin{cases} r S'_p(\psi_N, t) + \frac{1}{\mu_0 r} S_{ff'}(\psi_N, t) & \text{in } \Omega_P(\psi), \\ S_i^{-1} \left( \mathbf{S} \vec{V} + \mathbf{R} \vec{\Psi}(\partial_t \psi) \right)_i & \text{in } \Omega_{\text{coil}_i}, \\ -\frac{\sigma_k}{r} \partial_t \psi & \text{in } \Omega_{\text{passive}}, \\ 0 & \text{elsewhere,} \end{cases}$$

$$\psi(0, z, t) = 0, \quad \lim_{\|(r,z)\| \rightarrow +\infty} \psi(r, z, t) = 0, \quad \psi(r, z, 0) = \psi_0(r, z),$$

Circuit equations  $\vec{I} = \mathbf{S} \vec{V} + \mathbf{R} \vec{\Psi}(\partial_t \psi)$ :



where  $\vec{\Psi}(\partial_t \psi) = (\dots, \int_{\Omega_{\text{coil}_i}} \partial_t \psi dr dz, \dots)$ .

# Inverse Evolution Problem (for volt. evolution $\vec{V}(t)$ )

Objective and Regularization:

$$K(\psi(t)) := \frac{1}{2} \int_0^T \left( \sum_{i=1}^{N_{\text{desi}}} (\psi(r_i(t), z_i(t), t) - \psi(r_{\text{desi}}(t), z_{\text{desi}}(t), t))^2 \right) dt,$$

$$R(\vec{V}) := \sum_{i=1}^L \frac{w_i}{2} \int_0^T \vec{V}_i(t) \cdot \vec{V}_i(t) dt.$$

Optimal Control/Inverse Problem:

$$\min_{\psi(t), \vec{V}(t)} K(\psi(t)) + R(\vec{V})$$

subject to

$$-\nabla \cdot \left( \frac{1}{\mu r} \nabla \psi \right) = \begin{cases} r S'_p(\psi_N, t) + \frac{1}{\mu_0 r} S_{ff'}(\psi_N, t) & \text{in } \Omega_P(\psi), \\ S_i^{-1} \left( \mathbf{S} \vec{V} + \mathbf{R} \vec{\Psi}(\partial_t \psi) \right)_i & \text{in } \Omega_{\text{coil}_i}, \\ -\frac{\sigma_k}{r} \partial_t \psi & \text{in } \Omega_{\text{passive}}, \\ 0 & \text{elsewhere,} \end{cases}$$

$$\psi(0, z, t) = 0, \quad \lim_{\|(r,z)\| \rightarrow +\infty} \psi(r, z, t) = 0, \quad \psi(r, z, 0) = \psi_0(r, z),$$

# What's next?

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## Weak Formulation

Newton's Method

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# Weak Formulation

Find  $\psi \in V$  such that

$$A(\psi, \xi) - J_p(\psi, \xi) + c(\psi, \xi) = \ell(\vec{I}, \xi) \quad \forall \xi \in V.$$

with

$$V = \left\{ \psi : \Omega \rightarrow \mathbb{R}, \int_{\Omega} \psi^2 r \, dr dz < \infty, \int_{\Omega} (\nabla \psi)^2 r^{-1} \, dr dz < \infty, \psi|_{r=0} = 0 \right\},$$

$$A(\psi, \xi) := \int_{\Omega} \frac{1}{\mu r} \nabla \psi \cdot \nabla \xi \, dr dz, \quad \ell(\vec{I}, \xi) := \sum_{i=1}^{N_{\text{coil}}} S_i^{-1} \vec{I}_i \int_{\Omega_{\text{coil}_i}} \xi \, dr dz,$$

$$J_p(\psi, \xi) := \int_{\Omega_p(\psi)} \left( r p'(\psi) + \frac{1}{\mu_0 r} f f'(\psi) \right) \xi \, dr dz,$$

No integral equations in the spirit of FEM-BEM or "mariage à la mode"  
(Zienkiewics, Johnson, Nedelec, ...)

$c(\psi, \xi) \approx \int_{\partial\Omega} \xi \partial_n \psi \, dS$  taking into account boundary condition at infinity, details ...

What is  $\Omega$ , what is  $c(\cdot, \cdot)$ ?

## Weak Formulation, from infinite to finite

If  $\Omega$  semi-circle with radius  $\rho$  [Albanese1986, Gatica1995] ( $\Gamma = \partial\Omega$ ):

$$\begin{aligned}c(\psi, \xi) &= \int_{\partial\Omega} \xi \partial_n \psi dS = \int_{\Gamma} \xi(\mathbf{P}_1) \partial_{n_1} \left( \int_{\Gamma} \partial_{n_2} G(\mathbf{P}_1, \mathbf{P}_2) \psi(\mathbf{P}_2) dS_2 \right) dS_1 \\ &= \frac{1}{2\mu_0} \int_{\Gamma} \int_{\Gamma} \psi(\mathbf{P}_1) M(\mathbf{P}_1, \mathbf{P}_2) \xi(\mathbf{P}_2) dS_1 dS_2 \\ &= \frac{1}{\mu_0} \int_{\Gamma} \psi(\mathbf{P}_1) N(\mathbf{P}_1) \xi(\mathbf{P}_1) dS_1 \\ &\quad + \frac{1}{2\mu_0} \int_{\Gamma} \int_{\Gamma} (\psi(\mathbf{P}_1) - \psi(\mathbf{P}_2)) M(\mathbf{P}_1, \mathbf{P}_2) (\xi(\mathbf{P}_1) - \xi(\mathbf{P}_2)) dS_1 dS_2\end{aligned}$$

with  $G(\mathbf{P}_1, \mathbf{P}_2) \approx \log(\|\mathbf{P}_1 - \mathbf{P}_2\|)^{-1}$ , fundamental solution of  $\nabla \frac{1}{\mu_0 r} \nabla$  and

$$M(\mathbf{P}_1, \mathbf{P}_2) = \frac{k_{1,2}}{2\pi(r_1 r_2)^{\frac{3}{2}}} \left( \frac{2 - k_{1,2}^2}{2 - 2k_{1,2}^2} E(k_{1,2}) - K(k_{1,2}) \right)$$

where  $\mathbf{P}_i = (r_i, z_i)$ ,  $K$  and  $E$  complete elliptic integrals of 1st and 2nd kind, and

$$N(\mathbf{P}_1) = \int_{\Gamma} M(\mathbf{P}_1, \mathbf{P}_2) dS_2, \quad k_{1,2} = \sqrt{\frac{4r_1 r_2}{(r_1 + r_2)^2 + (z_1 - z_2)^2}}.$$

# Weak Formulation, evolution problem

Semi-discretization in time with implicit Euler

Set  $\psi^0 = \psi(t_0)$ . Find  $\psi^1, \dots, \psi^N \in V$  approximating  $\psi(t_1), \dots, \psi(t_N)$  with

$$\begin{aligned} \Delta t_k \mathbf{A}(\psi^k, \xi) - \Delta t_k \mathbf{J}_p^k(\psi^k, \xi) - j^{\text{ps}}(\psi^k, \xi) - j^{\text{c}}(\psi^k, \xi) + \Delta t_k c(\psi^k, \xi) \\ = \Delta t_k \ell(\mathbf{S}\vec{V}(t_k), \xi) - j^{\text{ps}}(\psi^{k-1}, \xi) - j^{\text{c}}(\psi^{k-1}, \xi) \quad \forall \xi \in V. \end{aligned}$$

$$j^{\text{ps}}(\psi, \xi) := - \sum_{i=1}^{N_{\text{passiv}}} \int_{\Omega_{\text{passive}}} \frac{\sigma_i}{r} \psi \xi \, dr dz, \quad j^{\text{c}}(\psi, \xi) := \sum_{i=1}^{N_{\text{coil}}} S_i^{-1} \left( \mathbf{R}\vec{\Psi}(\psi) \right)_i \int_{\Omega_{\text{coil}_i}} \xi \, dr dz.$$

Stationary Problem:

$$\mathbf{A}(\mathbf{y}) = \mathbf{F}(\mathbf{u}),$$

State  $\mathbf{y}$  is flux  $\psi$ ;  
Control  $\mathbf{u}$  are currents  $\vec{I}$ ;

Evolution Problem:

$$\mathbf{A}(\mathbf{y}^{k+1}) + \mathbf{m}(\mathbf{y}^{k+1}) = \mathbf{G}(\mathbf{u}^{k+1}) + \mathbf{m}(\hat{\mathbf{y}}^k),$$

State  $\mathbf{y}^1, \dots$  is flux  $\psi^1, \dots$ ;  
Control  $\mathbf{u}^1, \dots$  are voltages  $\vec{V}(t_1), \dots$ ;

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# Newton's Method, Continuous Approach

Nonlinear variational formulation

$$A(\psi, \xi) + J_p(\psi, \xi) + c(\psi, \xi) = \ell(\xi),$$

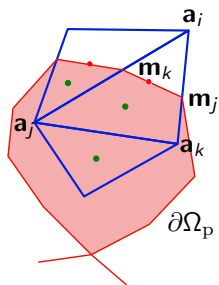
Newton iterations

$$(D_\psi A(\psi^k, \xi) + D_\psi J_p(\psi^k, \xi)) (\psi^{k+1} - \psi^k) = \ell(\xi) - A(\psi^k, \xi) - J_p(\psi^k, \xi) - c(\psi^k, \xi),$$

- ▶  $D_\psi A(\psi^k, \xi)$  simple,  $D_\psi J_p(\psi^k, \xi)$  not so simple;

From **shape derivatives**, **rearrangement**, or . . . .

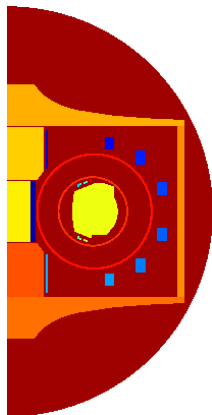
$$D_\psi J_p(\psi^k, \xi) \psi = \int_{\Omega_p(\psi^k)} D_\psi j_p(r, \psi^k(r, z)) \psi^k \xi_h dr dz + \int_{\partial\Omega_p(\psi^k)} j_p(r, \psi^k(r, z)) \xi \frac{\psi_{bd}(\psi^k) - \psi(r, z)}{|\nabla \psi^k|} dS,$$



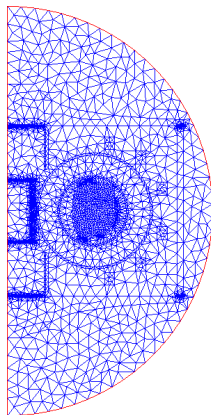
- ▶  $\psi_h$  is discretization of  $\psi$  with **linear finite elements**;
- ▶  $\partial\Omega_p(\psi_h^k)$  piecewise straight;
- ▶ barycenter quadrature rule for surface integrals
- ▶ midpoint quadrature rule for line integrals;

# Newton's Method: Example I

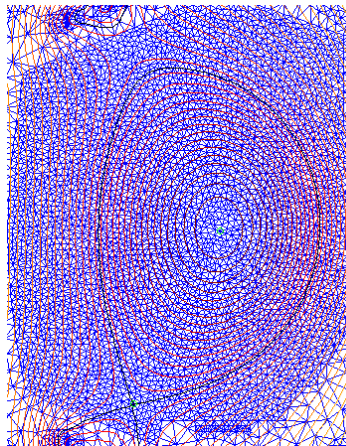
Subdomains



Mesh:



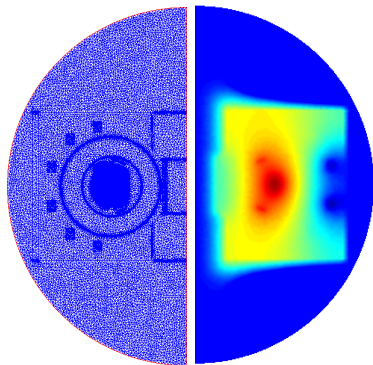
Solution:



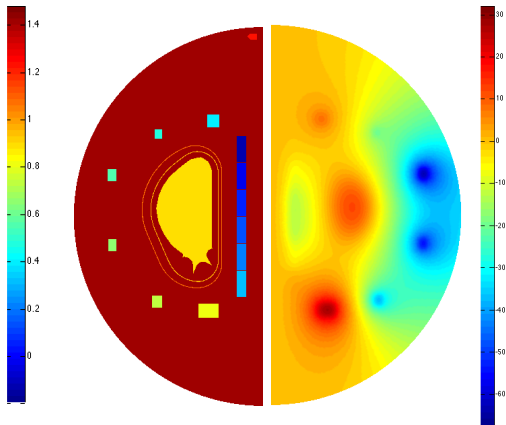
- ▶ mesh generation with TRIANGLE linear FEM;
- ▶ direct linear solver UMFPACK;

# Newton's Method: Example II

West (with Iron)

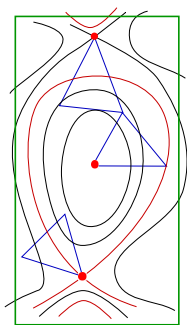


ITER (without Iron):

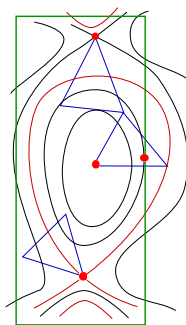
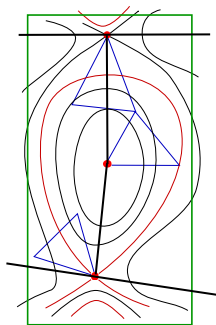


# Newton's Method, Plasma Domain $\Omega_p(\psi_h^k)$

How to determine plasma domain  $\Omega_p(\psi_h^k)$ : *Interior of last closed isoline*



X-point plasma



limiter plasma

Algorithm:

1. find the maximum location  $\mathbf{P}_{ax}$  of  $\psi_h$
2. find all **discrete** saddle points  $\mathbf{P}_X$  of  $\psi_h$ ;
3. construct excluded area, via perpendicular cut lines;
4.  $\psi_{bd}$  is the maximal value of all  $\psi_h(\mathbf{P}_X)$  and  $\psi_h$  on limiter;

# Newton's Method, Linearization I

## Newton iterations

$$(D_\psi \mathbf{A}(\psi^k, \xi) + D_\psi \mathbf{J}_p(\psi^k, \xi)) (\psi^{k+1} - \psi^k) = \ell(\xi) - \mathbf{A}(\psi^k, \xi) - \mathbf{J}_p(\psi^k, \xi) - c(\psi^k, \xi),$$

**Problem 1:** We did observe fast but *not quadratic* convergence!

**Problem 2:** Gradient test failed!

Recall Gradient test:

$$\begin{aligned} \widehat{C}(\mathbf{u}) &:= C(\mathbf{y}(\mathbf{u}), \mathbf{u}) & \Rightarrow & \nabla_{\mathbf{u}} \widehat{C}(\mathbf{u}) = \nabla_{\mathbf{u}} C(\mathbf{y}(\mathbf{u}), \mathbf{u}) + \nabla_{\mathbf{u}} \mathbf{F}(\mathbf{u})^T \mathbf{p} \\ \text{with } \mathbf{A}(\mathbf{y}(\mathbf{u})) &= \mathbf{F}(\mathbf{u}) & & \text{with } \nabla_{\mathbf{y}} \mathbf{A}(\mathbf{y}(\mathbf{u}))^T \mathbf{p} = -\nabla_{\mathbf{y}} C(\mathbf{y}(\mathbf{u}), \mathbf{u}) \\ & & & \left\| \frac{\widehat{C}(\mathbf{u} + \varepsilon \delta \mathbf{u}) - \widehat{C}(\mathbf{u})}{\varepsilon} - \nabla_{\mathbf{u}} \widehat{C}(\mathbf{u}) \delta \mathbf{u} \right\| = O(\varepsilon) \end{aligned}$$

**Reason:** Quadrature of analytic derivative  $D_\psi \mathbf{J}_p(\psi_h^k, \xi_h)$  is **not a derivative!**  
A discrete non-linear current:

$$\mathbf{J}_h(\psi_h, \xi_h) = \sum_T |T \cap \Omega_p(\psi_h)| j_p(\mathbf{b}_T, \psi_h(\mathbf{b}_T)) \xi_h(\mathbf{b}_T)$$

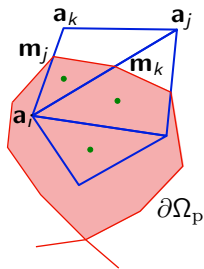
and  $\mathbf{b}_T = \mathbf{b}_T(\psi_h)$

# Newton's Method, Linearization II

A discrete non-linear current:

$$J_h(\psi_h, \xi_h) = \sum_T |T \cap \Omega_P(\psi_h)| j_P(\mathbf{b}_T, \psi_h(\mathbf{b}_T)) \xi_h(\mathbf{b}_T) \text{ and } \mathbf{b}_T = \mathbf{b}_T(\psi_h)$$

The derivative (the **true** discrete derivative!):



$$\begin{aligned} D_\psi J_h(\psi_h, \xi_h)(\lambda_i) &= \frac{d}{d\psi_i} J_h(\psi_h, \xi_h) \\ &= \sum_T \frac{d}{d\psi_i} |T \cap \Omega_P(\psi_h)| j_P(\mathbf{b}_T, \psi_N(\mathbf{b}_T)) \xi_h(\mathbf{b}_T) \\ &\quad + \sum_T |T \cap \Omega_P(\psi_h)| \frac{d}{d\psi_i} j_P(\mathbf{b}_T, \psi_N(\mathbf{b}_T)) \xi_h(\mathbf{b}_T) \\ &\quad + \sum_T |T \cap \Omega_P(\psi_h)| j_P(\mathbf{b}_T, \psi_N(\mathbf{b}_T)) \frac{d}{d\psi_i} \xi_h(\mathbf{b}_T) \end{aligned}$$

General implementation philosophy (everything local)

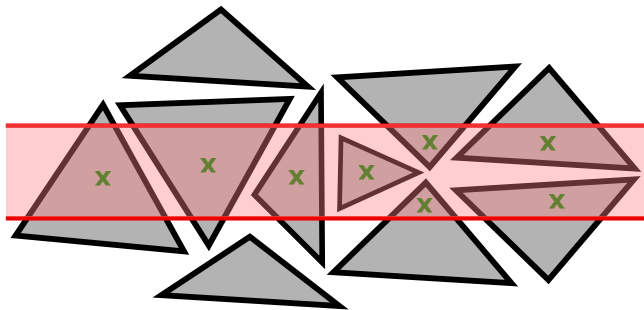
- ▶ Compute barycenter & intersection and **their derivatives** at same time!
- ▶ Assemble **vector**  $J_h(\psi_h, \lambda_j)$  and **matrix**  $D_\psi J_h(\psi_h, \lambda_j)(\lambda_i)$  at same time.

# Newton's Method; Levelset and Mesh

## Core function

Find intersection and quadrature points (and derivatives) of all elements that have non-zero intersection with levelline between  $\psi_l$  and  $\psi_u$ .

- ▶  $\psi_u = \psi_{ax}$  and  $\psi_l = \psi_{bnd}$  for plasma domain;



Centroid formula to generate quadrature formulas

$$Area_{tot} \mathbf{Bary}_{tot} = \sum_i Area_i \mathbf{Bary}_i$$

# Newton's Method, Linearization III

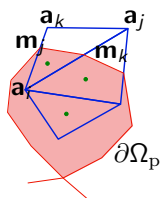
Case A:  $T \cap \Omega_p(\psi_h) = T$ :

$$\text{barycenter } (r_T, z_T) = \frac{1}{3}(\mathbf{a}_i + \mathbf{a}_j + \mathbf{a}_k).$$

Case B:  $T \cap \Omega_p(\psi_h) = \text{triangle}$

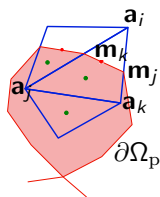
$$(r_T(\psi_h), z_T(\psi_h)) = \frac{1}{3}(\mathbf{a}_i + \mathbf{m}_k(\psi_h) + \mathbf{m}_j(\psi_h))$$

$$\mathbf{a}_i + \frac{1}{3}\lambda_j(\mathbf{m}_k(\psi_h))(\mathbf{a}_j - \mathbf{a}_i) + \frac{1}{3}\lambda_k(\mathbf{m}_j(\psi_h))(\mathbf{a}_k - \mathbf{a}_i)$$



Case C:  $T \cap \Omega_p(\psi_h) = \text{quadrilateral}$

$$(r_T(\psi_h), z_T(\psi_h)) = \mathbf{a}_i + \frac{1}{3} \frac{1 - \lambda_j^2(\mathbf{m}_k(\psi_h))\lambda_k(\mathbf{m}_j(\psi_h))}{1 - \lambda_j(\mathbf{m}_k(\psi_h))\lambda_k(\mathbf{m}_j(\psi_h))} (\mathbf{a}_j - \mathbf{a}_i) + \frac{1}{3} \frac{1 - \lambda_j(\mathbf{m}_k(\psi_h))\lambda_k^2(\mathbf{m}_j(\psi_h))}{1 - \lambda_j(\mathbf{m}_k(\psi_h))\lambda_k(\mathbf{m}_j(\psi_h))} (\mathbf{a}_k - \mathbf{a}_i)$$



$$\lambda_j(\mathbf{m}_k(\psi_h)) = \frac{\psi_{\text{bd}}(\psi_h) - \psi_i}{\psi_j - \psi_i}, \quad \lambda_k(\mathbf{m}_j(\psi_h)) = \frac{\psi_{\text{bd}}(\psi_h) - \psi_i}{\psi_k - \psi_i},$$



# Newton's Method, "Semi-Automatic" Differentiation

```
function [J_NL,DJ] = assemblePlasma_NL(Mesh,psi,jpHandle)
% J_NL: vector, non-linear operator at psi,
% DJ: matrix, derivative of J_NL at psi
...
% levelset is structure levelset.ratio, level.bary
% containing ratio of intersection domain, barycenter
% and derivatives for each element
levelset = find_Plasma(Mesh,psi);
....
....
% non-linear operator
J_NL = [0.5*det_BK.*ratio(:,1).*jplasma_bary(:,1).*N1(:,1);...
        0.5*det_BK.*ratio(:,1).*jplasma_bary(:,1).*N2(:,1);...
        0.5*det_BK.*ratio(:,1).*jplasma_bary(:,1).*N3(:,1)]
....
....
% derivative of non-linear operator
DJ.E = [0.5*det_BK.*ratio(:,2).*jplasma_bary(:,1).*N1(:,1)+...
        0.5*det_BK.*ratio(:,1).*jplasma_bary(:,2).*N1(:,1)+...
        0.5*det_BK.*ratio(:,1).*jplasma_bary(:,1).*N1(:,2);
        0.5*det_BK.*ratio(:,2).*jplasma_bary(:,1).*N2(:,1)+...
        0.5*det_BK.*ratio(:,1).*jplasma_bary(:,2).*N2(:,1)+...
        0.5*det_BK.*ratio(:,1).*jplasma_bary(:,1).*N2(:,2);
        0.5*det_BK.*ratio(:,2).*jplasma_bary(:,1).*N3(:,1)+...
        0.5*det_BK.*ratio(:,1).*jplasma_bary(:,2).*N3(:,1)+...
        0.5*det_BK.*ratio(:,1).*jplasma_bary(:,1).*N3(:,2);
```

# What's next?

Quasi-Static Free-Boundary Equilibrium of Toroidal Plasma

Direct Static Problem

Inverse Static Problem

Direct Evolution Problem

Inverse Evolution Problem

Weak Formulation

Newton's Method

**Sequential Quadratic Programming**

Validation & Performance

Application: Vertical Displacement

Conclusions & Outlook

# Inverse Static Problem (for currents $I_j$ )

Objective and Regularization

$$K(\psi) := \frac{1}{2} \sum_{i=1}^{N_{\text{desi}}} (\psi(r_i, z_i) - \psi(r_{\text{desi}}, z_{\text{desi}}))^2, \quad R(I_1, \dots, I_L) := \sum_{i=1}^L \frac{w_i}{2} I_i^2$$

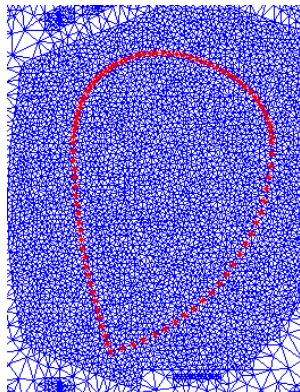
Optimal Control/Inverse Problem:

$$\min_{\psi, I_1, \dots, I_L} K(\psi) + R(I_1, \dots, I_L)$$

subject to

$$-\nabla \cdot \left( \frac{1}{\mu r} \nabla \psi \right) = \begin{cases} r S'_p(\psi_N) + \frac{1}{\mu_0 r} S_{ff'}(\psi_N) & \text{in } \Omega_P(\psi), \\ \frac{I_{i,j}}{S_{i,j}} & \text{in } \Omega_{\text{coil}_i}, \\ 0 & \text{elsewhere,} \end{cases}$$

$$\psi(0, z) = 0, \quad \lim_{\|(r,z)\| \rightarrow +\infty} \psi(r, z) = 0,$$



PDE-constrained optimization with **non-linear** constraints.

# Optimal Control: Numerical Methods

1.) unconstrained optimization:

$$\min_{\mathbf{u}} C(\mathbf{u}) \iff \nabla_{\mathbf{u}} C(\mathbf{u}) = 0$$

$$\underbrace{\mathbf{u}^{k+1} = \mathbf{u}^k - s_k \nabla C(\mathbf{u}^k)}_{\text{global, slow convergence}} \quad \underbrace{\mathbf{u}^{k+1} = \mathbf{u}^k - \mathbf{M}_k \nabla C(\mathbf{u}^k)}_{\text{global, not too slow convergence}} \quad \underbrace{\mathbf{u}^{k+1} = \mathbf{u}^k - \nabla^2 C(\mathbf{u}^k)^{-1} \nabla C(\mathbf{u}^k)}_{\text{local, fast convergence}}$$

2.) constrained optimization:

$$\begin{aligned} & \min_{\mathbf{y}, \mathbf{u}} C(\mathbf{y}, \mathbf{u}) \\ & \text{s.t.} \quad \mathbf{A}_1(\mathbf{y}) = \mathbf{F}_1(\mathbf{u}) \\ & \quad \quad \quad \vdots \\ & \quad \quad \quad \mathbf{A}_n(\mathbf{y}) = \mathbf{F}_n(\mathbf{u}) \end{aligned} \iff (\text{A}) \iff (\text{B})$$

(A) reduced approach/steepest descent:

$$\min_{\mathbf{u}} \hat{C}(\mathbf{u}) := C(\mathbf{y}(\mathbf{u}), \mathbf{u}) \iff \begin{aligned} \text{i) } & \nabla_{\mathbf{u}} \hat{C}(\mathbf{u}) = \nabla_{\mathbf{u}} C(\mathbf{y}(\mathbf{u}), \mathbf{u}) + \nabla_{\mathbf{y}} C(\mathbf{y}(\mathbf{u}), \mathbf{u}) \nabla_{\mathbf{u}} \mathbf{y}(\mathbf{u}) = 0 \\ & \text{with } \nabla_{\mathbf{y}} \mathbf{A}(\mathbf{y}(\mathbf{u})) \nabla_{\mathbf{u}} \mathbf{y}(\mathbf{u}) = \nabla_{\mathbf{u}} \mathbf{F}(\mathbf{u}) \\ \text{ii) } & \nabla_{\mathbf{u}} \hat{C}(\mathbf{u}) = \nabla_{\mathbf{u}} C(\mathbf{y}(\mathbf{u}), \mathbf{u}) + \nabla_{\mathbf{u}} \mathbf{F}(\mathbf{u})^T \mathbf{p} = 0 \\ & \text{with } \nabla_{\mathbf{y}} \mathbf{A}(\mathbf{y}(\mathbf{u}))^T \mathbf{p} = -\nabla_{\mathbf{y}} C(\mathbf{y}(\mathbf{u}), \mathbf{u}) \end{aligned}$$

To use methods from 1) we need to compute  $\mathbf{y}(\mathbf{u})$  and  $\nabla_{\mathbf{u}} \mathbf{y}(\mathbf{u})$ ! expensive!

(B) Lagrange multipliers  $p_i$ / SQP:

stationary point of Lagrangian  $L(\mathbf{y}, \mathbf{u}, \mathbf{p}) = C(\mathbf{y}, \mathbf{u}) + \mathbf{p}^T (\mathbf{A}(\mathbf{y}) - \mathbf{F}(\mathbf{u}))$

**Fastest algorithm!** Newton for (B) = Sequential quadratic programming (SQP)

# Steepest Descent vs. SQP I: fminunc

## Command Window

Iteration	f(x)	step	optimality	CG-iteration
0	0.0748239		3.33e-06	
1	0.0747357	10	3.33e-06	2
2	0.0745595	20	3.33e-06	2
3	0.0742077	40	3.32e-06	2
4	0.073506	80	3.32e-06	2
5	0.0721101	160	3.3e-06	2
6	0.0693473	320	3.27e-06	2
7	0.0644181	640	2.55e-06	2
8	0.0562001	1280	2.44e-06	1
16	0.00655371	40960	6.72e-07	4
17	0.00603318	669.542	3.49e-08	1
18	0.00603318	81920	3.49e-08	7
19	0.00603318	20480	3.49e-08	0
20	0.00603318	5120	3.49e-08	0
21	0.00600083	1280	6.15e-08	0
22	0.00592413	2560	3.2e-08	5
23	0.00581343	5120	5.59e-08	7
24	0.00581343	10240	5.59e-08	5
25	0.00581343	2560	5.59e-08	0
26	0.00581343	640	5.59e-08	0
27	0.00581343	160	5.59e-08	0
28	0.00581343	40	5.59e-08	0
29	0.00581224	10	5.16e-08	0
30	0.0058102	20	4.32e-08	5

Solver stopped prematurely.

fminunc stopped because it exceeded the function evaluation limit, options.MaxFunEvals = 30 (the selected value).

Elapsed time is 206.081131 seconds.

# Steepest Descent vs. SQP II: home made sqp

## Command Window

```
Newtoniteration 0; relativ residuum 3.84e-05 ←  
cost 7.41e-02; regularization 7.25e-04; objective 7.48e-02  
resid_stat = 2.00e+00, resid_ctrl = 2.31e+03, resid_adj = 1.13e+00  
Newtoniteration 1; relativ residuum 5.86e-01 ←  
cost 1.70e-03; regularization 1.50e-03; objective 3.20e-03  
resid_stat = 5.41e-01, resid_ctrl = 5.47e-13, resid_adj = 5.86e-01  
Newtoniteration 2; relativ residuum 2.84e-01 ←  
cost 1.90e-03; regularization 1.34e-03; objective 3.24e-03  
resid_stat = 5.90e-01, resid_ctrl = 2.21e-16, resid_adj = 2.84e-01  
Newtoniteration 3; relativ residuum 4.34e-02 ←  
cost 1.86e-03; regularization 1.38e-03; objective 3.24e-03  
resid_stat = 1.48e-03, resid_ctrl = 2.20e-16, resid_adj = 4.34e-02  
Newtoniteration 4; relativ residuum 4.50e-03 ←  
cost 1.86e-03; regularization 1.38e-03; objective 3.24e-03  
resid_stat = 2.92e-04, resid_ctrl = 5.05e-16, resid_adj = 4.50e-03  
Newtoniteration 5; relativ residuum 1.70e-04 ←  
cost 1.86e-03; regularization 1.38e-03; objective 3.24e-03  
resid_stat = 1.23e-04, resid_ctrl = 5.30e-16, resid_adj = 1.70e-04  
Newtoniteration 6; relativ residuum 1.54e-05 ←  
cost 1.86e-03; regularization 1.38e-03; objective 3.24e-03  
resid_stat = 5.67e-05, resid_ctrl = 2.65e-16, resid_adj = 1.54e-05  
Newtoniteration 7; relativ residuum 4.06e-10 ←  
cost 1.86e-03; regularization 1.38e-03; objective 3.24e-03  
resid_stat = 7.50e-07, resid_ctrl = 4.21e-16, resid_adj = 4.06e-10  
Newtoniteration 8; relativ residuum 1.12e-12 ←  
cost 1.86e-03; regularization 1.38e-03; objective 3.24e-03  
resid_stat = 4.01e-09, resid_ctrl = 3.16e-16, resid_adj = 1.12e-12
```

Elapsed time is 11.649052 seconds.

# Inverse Evolution Problem (for volt. evolution $\vec{V}(t)$ )

Objective and Regularization:

$$K(\psi(t)) := \frac{1}{2} \int_0^T \left( \sum_{i=1}^{N_{\text{desi}}} (\psi(r_i(t), z_i(t), t) - \psi(r_{\text{desi}}(t), z_{\text{desi}}(t), t))^2 \right) dt,$$

$$R(\vec{V}) := \sum_{i=1}^L \frac{w_i}{2} \int_0^T \vec{V}_i(t) \cdot \vec{V}_i(t) dt.$$

Optimal Control/Inverse Problem:

$$\begin{aligned} & \min_{\psi(t), \vec{V}(t)} K(\psi(t)) + R(\vec{V}) \\ & \text{subject to} \\ & -\nabla \cdot \left( \frac{1}{\mu r} \nabla \psi \right) = \begin{cases} r S_p'(\psi_N, t) + \frac{1}{\mu_0 r} S_{\#}'(\psi_N, t) & \text{in } \Omega_P(\psi), \\ S_i^{-1} \left( \mathbf{S} \vec{V} + \mathbf{R} \dot{\Psi}(\partial_t \psi) \right)_i & \text{in } \Omega_{\text{coil}_i}, \\ -\frac{\sigma_k}{r} \partial_t \psi & \text{in } \Omega_{\text{passive}}, \\ 0 & \text{elsewhere,} \end{cases} \\ & \psi(0, z, t) = 0, \quad \lim_{\|(r,z)\| \rightarrow +\infty} \psi(r, z, t) = 0, \quad \psi(r, z, 0) = \psi_0(r, z), \end{aligned}$$

PDE-constrained optimization with **non-linear** constraints.





# Inverse Evolution Problem, Details

$$D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}}) = \begin{pmatrix} \tau D_{\mathbf{y}}\mathbf{A}(\mathbf{y}_1) + D_{\mathbf{y}}\mathbf{m}(\mathbf{y}_1) & 0 & 0 \\ -D_{\mathbf{y}}\mathbf{m}(\mathbf{y}_1) & \tau D_{\mathbf{y}}\mathbf{A}(\mathbf{y}_2) + D_{\mathbf{y}}\mathbf{m}(\mathbf{y}_2) & 0 \\ \vdots & \vdots & \vdots \\ 0 & -D_{\mathbf{y}}\mathbf{m}(\mathbf{y}_{N-1}) & \tau D_{\mathbf{y}}\mathbf{A}(\mathbf{y}_N) + D_{\mathbf{y}}\mathbf{m}(\mathbf{y}_N) \end{pmatrix},$$

$$D_{\bar{\mathbf{u}}}\bar{\mathbf{F}}(\bar{\mathbf{u}}) = \begin{pmatrix} \tau D_{\mathbf{u}}\tilde{\mathbf{F}}(\mathbf{u}_1) & 0 & 0 \\ 0 & \tau D_{\mathbf{u}}\tilde{\mathbf{F}}(\mathbf{u}_2) & 0 \\ \vdots & \vdots & \vdots \\ 0 & 0 & \tau D_{\mathbf{u}}\tilde{\mathbf{F}}(\mathbf{u}_N) \end{pmatrix}, \bar{\mathbf{K}} = \begin{pmatrix} \tau \mathbf{K}_1 & 0 & 0 \\ 0 & \tau \mathbf{K}_2 & 0 \\ \vdots & \vdots & \vdots \\ 0 & 0 & \tau \mathbf{K}_N \end{pmatrix},$$

$$\bar{\mathbf{R}} = \begin{pmatrix} \tau \mathbf{R}_1 & 0 & 0 & 0 \\ 0 & \tau \mathbf{R}_2 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \tau \mathbf{R}_N \end{pmatrix},$$

# Inverse Evolution Problem

## Sequential Quadratic Programming

"Solve a sequence of quadratic problems"

1.) Solve full Newton system:

Quasi-Newton for  $\bar{\mathbf{y}}^{k+1} = \bar{\mathbf{y}}^k + \Delta\bar{\mathbf{y}}$ ,  $\bar{\mathbf{u}}^{k+1} = \bar{\mathbf{u}}^k + \Delta\bar{\mathbf{u}}$ ,  $\bar{\mathbf{p}}^{k+1} = \bar{\mathbf{p}}^k + \Delta\bar{\mathbf{p}}$ :

$$\begin{pmatrix} \bar{\mathbf{K}} & 0 & D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}^T(\bar{\mathbf{y}}^k) \\ 0 & \bar{\mathbf{R}} & -D_{\bar{\mathbf{u}}}\bar{\mathbf{F}}^T(\bar{\mathbf{u}}^k) \\ D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}}^k) & -D_{\bar{\mathbf{u}}}\bar{\mathbf{F}}(\bar{\mathbf{u}}^k) & 0 \end{pmatrix} \begin{pmatrix} \Delta\bar{\mathbf{y}} \\ \Delta\bar{\mathbf{u}} \\ \Delta\bar{\mathbf{p}} \end{pmatrix} = - \begin{pmatrix} \bar{\mathbf{K}}\bar{\mathbf{y}}^k + D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}^T(\bar{\mathbf{y}}^k)\bar{\mathbf{p}}^k \\ \bar{\mathbf{R}}\bar{\mathbf{u}}^k - D_{\bar{\mathbf{u}}}\bar{\mathbf{F}}^T(\bar{\mathbf{u}}^k)\bar{\mathbf{p}}^k \\ \bar{\mathbf{A}}(\bar{\mathbf{y}}^k) - \bar{\mathbf{A}}_0(\mathbf{y}_0) - \bar{\mathbf{F}}(\bar{\mathbf{u}}^k) \end{pmatrix},$$

System is **roughly twice as large** as for direct problem.

2.) Solve reduced Newton system with CG or directly:

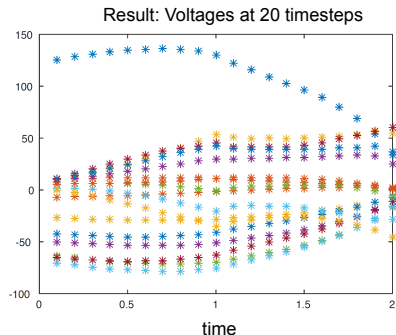
Eliminate  $\Delta\bar{\mathbf{y}}$ ,  $\Delta\bar{\mathbf{p}}$ :  $\bar{\mathbf{M}}(\bar{\mathbf{u}}^k, \bar{\mathbf{y}}^k)\Delta\bar{\mathbf{u}} = \bar{\mathbf{h}}(\bar{\mathbf{u}}^k, \bar{\mathbf{y}}^k)$ ,

where  $\bar{\mathbf{M}}(\bar{\mathbf{u}}, \bar{\mathbf{y}}) = \bar{\mathbf{R}} + D_{\bar{\mathbf{u}}}\bar{\mathbf{F}}^T(\bar{\mathbf{u}})D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}})^{-T}\bar{\mathbf{K}}D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}})^{-1}D_{\bar{\mathbf{u}}}\bar{\mathbf{F}}(\bar{\mathbf{u}}^k)$

and  $\bar{\mathbf{h}}(\bar{\mathbf{u}}, \bar{\mathbf{y}}) = -\bar{\mathbf{R}}\bar{\mathbf{u}} - D_{\bar{\mathbf{u}}}\bar{\mathbf{F}}^T(\bar{\mathbf{u}})D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}})^{-T}\bar{\mathbf{K}}D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}})^{-1}(\bar{\mathbf{F}}(\bar{\mathbf{u}}) - \bar{\mathbf{A}}(\bar{\mathbf{y}}) + D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}})\bar{\mathbf{y}})$ .

- ▶ CG-solver: **Very few** CG-iterations, but one inversion of  $D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}})$  and  $D_{\bar{\mathbf{y}}}\bar{\mathbf{A}}(\bar{\mathbf{y}})^T$  in each iteration.
- ▶ direct solver:  $\bar{\mathbf{M}}(\bar{\mathbf{u}}^k, \bar{\mathbf{y}}^k)$  is relatively small but **not sparse**.

# Control of Transient Plasma Equilibrium



Major critics: electrodynamic effects of the plasma are not included, yet!

Remaining equations:

- ▶ conservation of density and energy;
- ▶  $-\partial_t \mathbf{B}_T = \mathbf{curl} \mathbf{E}_P$  in plasma;
- ▶  $\mathbf{E} + \mathbf{v} \times \mathbf{B} = \eta \mathbf{J}$  in plasma;

# What's next?

Quasi-Static Free-Boundary Equilibrium of Toroidal Plasma

Direct Static Problem

Inverse Static Problem

Direct Evolution Problem

Inverse Evolution Problem

Weak Formulation

Newton's Method

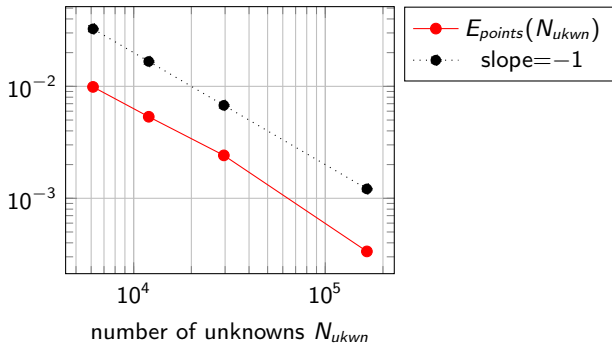
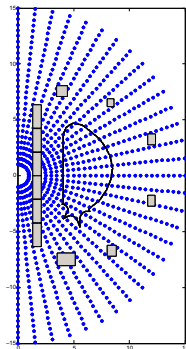
Sequential Quadratic Programming

**Validation & Performance**

Application: Vertical Displacement

Conclusions & Outlook

# Validation & Performance



number of triangles	number of unknowns	computing time (in s)	iteration	relative residual
12099	6134	2	1	$2.667473 \times 10^{+00}$
23723	11985	5	2	$9.157459 \times 10^{-02}$
58744	29556	11	3	$1.781645 \times 10^{-03}$
328693	164887	88	4	$0.525234 \times 10^{-06}$
1153174	577415	368	5	$3.935226 \times 10^{-12}$

# What's next?

Quasi-Static Free-Boundary Equilibrium of Toroidal Plasma

Direct Static Problem

Inverse Static Problem

Direct Evolution Problem

Inverse Evolution Problem

Weak Formulation

Newton's Method

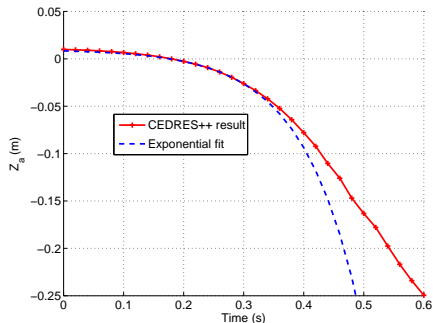
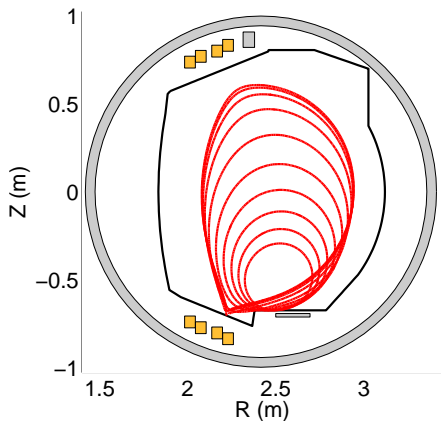
Sequential Quadratic Programming

Validation & Performance

**Application: Vertical Displacement**

Conclusions & Outlook

## Application: Vertical Displacement



**Figure:** Left: Plasma boundary at intervals of 100ms in a vertical instability simulation for WEST. Right: Time evolution of the vertical position of the magnetic axis  $z_{ax}$  in a vertical instability simulation for WEST.

# What's next?

Quasi-Static Free-Boundary Equilibrium of Toroidal Plasma

Direct Static Problem

Inverse Static Problem

Direct Evolution Problem

Inverse Evolution Problem

Weak Formulation

Newton's Method

Sequential Quadratic Programming

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Application: Vertical Displacement

Conclusions & Outlook



# Conclusions & Outlook

## Conclusions:

- ▶ mature and sound equilibrium calculation;
- ▶ ready to use for **applications** and **automation**;
- ▶ Coupling of CEDRES and ETS (European Transport Solver) in ITM; (C. Boulbe & B. Faugas with J.F. Artaud, P. Huyn, V. Basiuk, E. Nardon, J. Urban, D. Kalupin at CEA, Munich, Prag)
- ▶ FEEQS.M with Edge Plasma Code for divertor load optimization; (H.H. with M. Bloomart, T. Baelmans, N. Gauger, D. Reiter at Jülich, Leuven, Kaiserslautern),

## Outlook:

1. towards monolithic solver for equilibrium and transport;
2. optimal control for scenario optimization for tokamaks;
3. control engineers are interested in realtime solutions of the coupled problem;
4. can not be achieved by only increasing computational power;